

SIMC L/T/LT/TT parameterization

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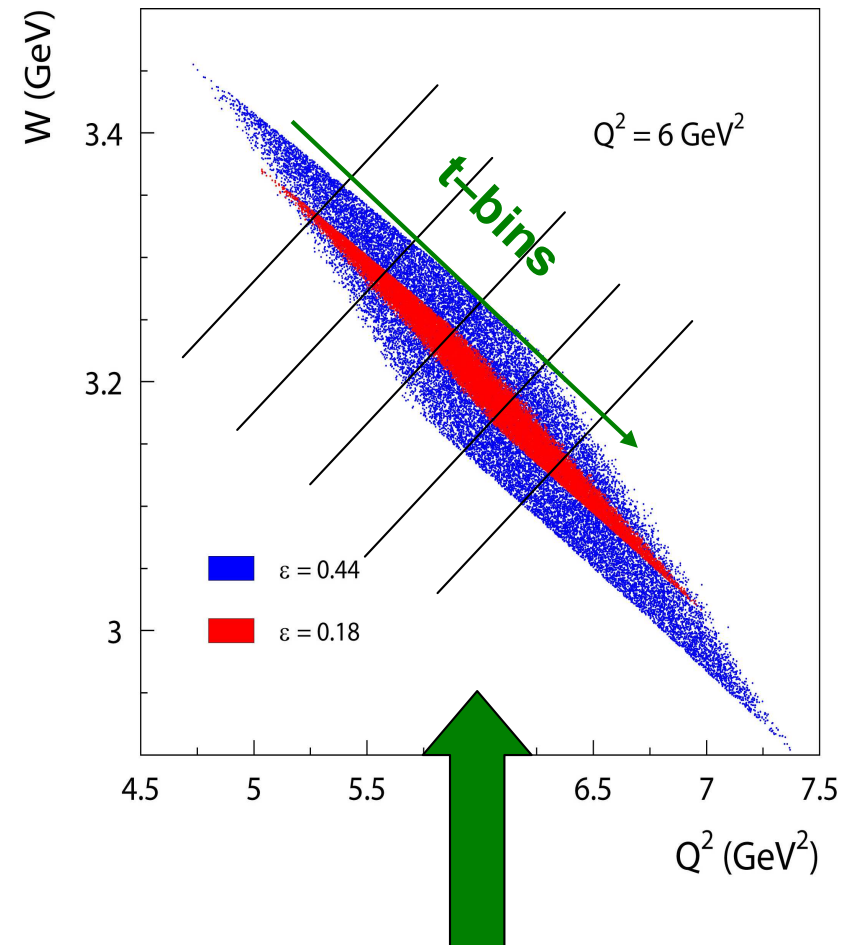
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Pre-requisite: Stable and finalized data!

- Before starting the L/T/LT/TT-separation procedure, it is essential that you have:
 - Final normalized yields (counts/mC) for all settings, with all efficiencies, livetimes, cryotarget, FADC-DT and other yield corrections tested for reliability over a wide rate range and applied
 - All kinematic offsets determined and finalized
- This is because it is essential that the one thing that is kept constant in the iterations is the experimental normalized yield and distributions.
- If any subsequent changes are made to any part of the experimental distributions, the iteration procedure must be repeated to ensure the result remains self-consistent.
- Failure to respect this restriction will result in significant wasted time!

Pre-requisite – Choose a functional form

- The cross section varies across experimental acceptance.
- It is needed to choose a functional form that will reasonably take into account this variation.
- Of course you don't know in advance what to choose, therein the uncertainty.
- All you can do is to make a choice, and start the iteration process with it.
- Then you need to do tests to see if the functional dependence reproduces the variation of the data.
- If the tests fail, then you need to modify the functional form and try again until you get something that works better.



Each t bin has a different average value of W , Q^2 .
This dependence must be taken into account

Example Parameterizations Used

- Fpi-1 LH+
- Blok et al., PRC 78 (2008) 045202

This procedure was carried out independently for Fpi-1 and Fpi-2 in order to have optimal descriptions in the two different kinematic ranges covered.⁴ The final cross section parametrization for Fpi-1 (the cross sections have units of $\mu\text{b}/\text{GeV}^2$, and the units of Q^2 , t , and m_π^2 are GeV^2) is

$$\begin{aligned}\frac{d\sigma_L}{dt} &= 36.51 e^{(26.10 - 7.75 Q^2)(t + 0.02)}, \\ \frac{d\sigma_T}{dt} &= \frac{0.74}{Q^2} + \frac{1.25}{Q^4} + 0.57 \frac{|t|}{(|t| + m_\pi^2)^2}, \\ \frac{d\sigma_{LT}}{dt} &= \left(\exp \left[4.69 + \frac{24.55}{\sqrt{Q^2}} t \right] + 1.47 - \frac{7.89}{Q^4} \right) \sin \theta^*, \\ \frac{d\sigma_{TT}}{dt} &= \left(\frac{3.44}{Q^2} - \frac{7.57}{Q^4} \right) \frac{|t|}{(|t| + m_\pi^2)^2} \sin^2 \theta^*. \quad (15)\end{aligned}$$

This parametrization is valid in the Q^2 range between 0.4 and 1.8 GeV^2 .

- Fpi-2 LH+
- Blok et al., PRC 78 (2008) 045202

The Fpi-2 parametrization, valid between $Q^2 = 1.4$ and 2.7 GeV^2 , is

$$\begin{aligned}\frac{d\sigma_L}{dt} &= \frac{350 Q^2}{(1 + 1.77 Q^2 + 0.05 Q^4)^2} e^{(16 - 7.5 \ln Q^2)t}, \\ \frac{d\sigma_T}{dt} &= \frac{4.5}{Q^2} + \frac{2.0}{Q^4}, \\ \frac{d\sigma_{LT}}{dt} &= \left(\exp \left[0.79 + \frac{3.4}{\sqrt{Q^2}} t \right] + 1.1 - \frac{3.6}{Q^4} \right) \sin \theta^*, \\ \frac{d\sigma_{TT}}{dt} &= -\frac{5.0}{Q^4} \frac{|t|}{(|t| + m_\pi^2)^2} \sin^2 \theta^*. \quad (16)\end{aligned}$$

It's okay to have different parameterizations for different Q^2 , W ranges, if that's what's required to have a good fit

Example Parameterizations Used

- Fpi-1 LD+
- Huber et al., PRC **91** (2015) 015202

This procedure was carried out independently for π^+ and π^- at each Q^2 , to have optimal descriptions in the different kinematic ranges covered. The parameterizations used in the $F_{\pi-1} \pi^+$ analysis are

$$\begin{aligned}\frac{d\sigma_L}{dt} &= g(W)[p_1 + p_2 \ln(Q^2)]e^{[p_3 + p_4 \ln(Q^2)](-t)}, \\ \frac{d\sigma_T}{dt} &= g(W)\left(\frac{|t| - |t_{\text{ave}}|}{|t_{\text{ave}}|}\right)\{p_5 + p_6 \ln(Q^2) \\ &\quad + [p_7 + p_8 \ln(Q^2)]\}, \\ \frac{d\sigma_{LT}}{dt} &= g(W)p_9 e^{p_{10}(-t)} \sin \theta_{\text{c.m.}}, \\ \frac{d\sigma_{TT}}{dt} &= g(W)f(t)\frac{p_{11}}{Q^2}e^{-Q^2} \sin^2 \theta_{\text{c.m.}},\end{aligned}\quad (9)$$

where $g(W) = 1/(W^2 - m_p^2)^2$ is the assumed W dependence discussed earlier, $f(t) = -t/(-t - m_\pi^2)^2$ is the pion pole factor, $|t_{\text{ave}}|$ is the average $-t$ value for a given kinematic setting, given by $|t_{\text{ave}}| = [0.105 + 0.04 \ln(Q^2)]Q^2$, and $p_{i=1,\dots,12}$ are the fit parameters.

- Fpi-1 LD-
- Huber et al., PRC **91** (2008) 015202

For the $F_{\pi-1} \pi^-$ analysis, a slightly different parametrization (because σ_T and σ_{TT} showed a stronger Q^2 -dependence) yielded a better fit:

$$\begin{aligned}\frac{d\sigma_L}{dt} &= g(W)[p_1 + p_2 \ln(Q^2)]e^{[p_3 + p_4 \ln(Q^2)](-t)}, \\ \frac{d\sigma_T}{dt} &= g(W)\left\{p_5 + \frac{p_6}{Q^4 + 0.1} \right. \\ &\quad \left. + [p_7 + p_8 \ln(Q^2)]\left(\frac{|t| - |t_{\text{ave}}|}{|t_{\text{ave}}|}\right)\right\}, \\ \frac{d\sigma_{LT}}{dt} &= g(W)p_9 e^{p_{10}(-t)} \sin \theta_{\text{c.m.}}, \\ \frac{d\sigma_{TT}}{dt} &= g(W)f(t)\left(\frac{p_{11}}{Q^2} + \frac{p_{12}}{Q^4 + 0.2}\right) \sin^2 \theta_{\text{c.m.}}.\end{aligned}\quad (10)$$

Pay attention to note above that LD-T, TT Q^2 -dependences had to be modified to yield a better fit

Example Parameterizations Used

- Fpi-2 LD+/-
- Huber et al., PRC **91** (2015) 015202

In the $F_{\pi-2}$ analyses, a common parametrization (similar to those in $F_{\pi-1}$) was used for both π^+ and π^- ,

$$\begin{aligned}\frac{d\sigma_L}{dt} &= g(W)[p_1 + p_2 \ln(Q^2)]e^{[p_3 + p_4 \ln(Q^2)](-t-0.2)}, \\ \frac{d\sigma_T}{dt} &= g(W)\left\{ p_5 + p_6 \ln(Q^2) \right. \\ &\quad \left. + [p_7 + p_8 \ln(Q^2)]\left(\frac{|t| - |t_{ave}|}{|t_{ave}|}\right) \right\}, \\ \frac{d\sigma_{LT}}{dt} &= g(W)\left[p_9 e^{p_{10}(-t)} + \frac{p_{11}}{(-t)} \right] \sin \theta_{c.m.}, \\ \frac{d\sigma_{TT}}{dt} &= g(W)f(t)\frac{p_{12}}{Q^2}e^{-Q^2} \sin^2 \theta_{c.m.},\end{aligned}\tag{11}$$

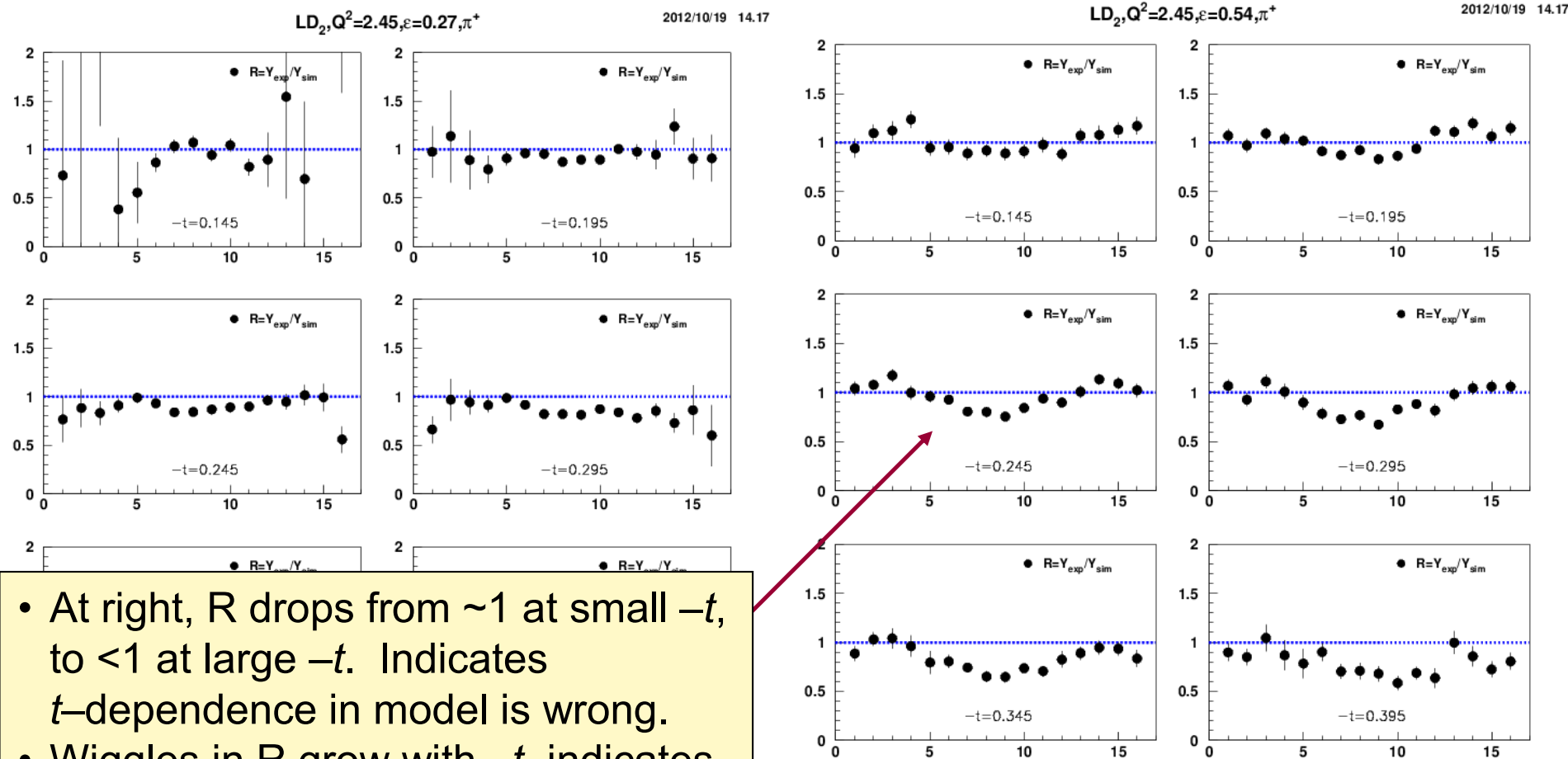
where $|t_{ave}| = [0.0735 + 0.028 \ln(Q^2)]Q^2$ and $p_4 = 0$.

LT and TT have extra $\sin\theta^*$ dependences since they are required to vanish in parallel-kinematics limit, i.e. $\theta^*=0$

The $|t|-|t_{ave}|$ factors vanish near the mean t for each (Q^2, W) kinematics. We wanted to incorporate a t -dependence that would only be a small correction away from t_{ave} , to make the iterations converge more easily. You would need your own calculation of t_{ave}

Evaluating if Fit parameters are okay

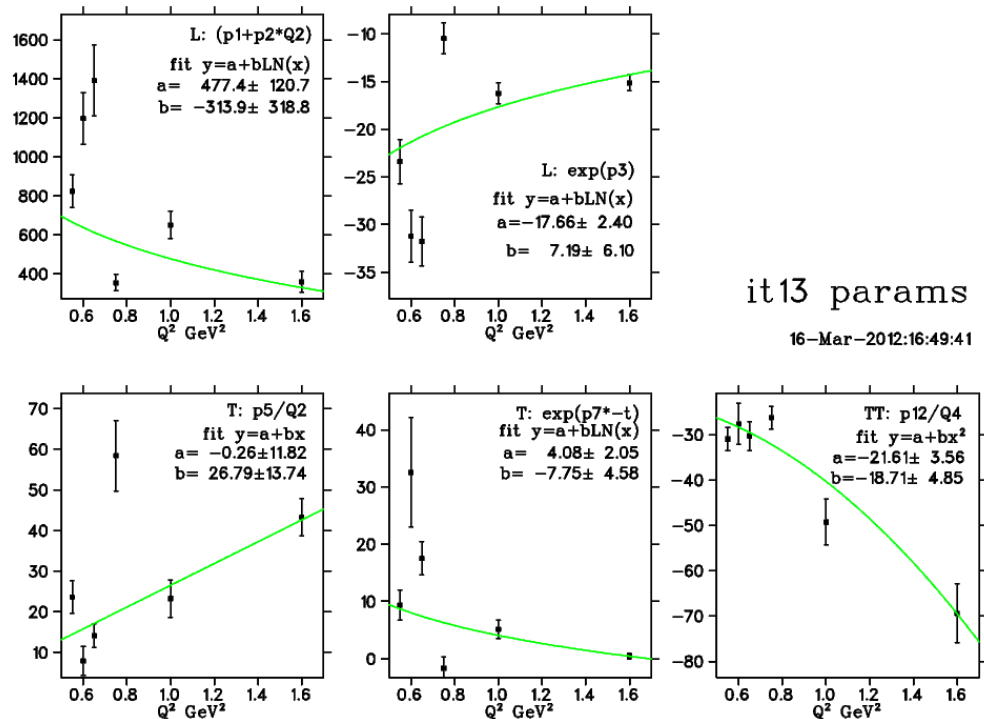
- Deviations between Data and MC usually are indicated as wiggles in R . We desire $R \approx 1$ over broad kinematic range



- At right, R drops from ~ 1 at small $-t$, to < 1 at large $-t$. Indicates t -dependence in model is wrong.
- Wiggles in R grow with $-t$, indicates problem with interference terms

Evaluating if Fit equations are okay

- Fit parameters are not required to be the same for different kinematics, but a consistent trend with Q^2 could indicate the parameterization equations need adjustment
- We compared `fitpar()` from different Q^2, W to see if they were slowly varying
- Here, $p5/Q^2$ parameter in T and $p12/Q^4$ parameter in TT grow with Q^2 , indicating a steeper Q^2 -dependence could be needed in model



More details are found in LD2
analysis technical report
HallC-docDB-773

SIMC modifications

- Replace physics_pion.f or physics_kaon.f with physics_iterate.f in SIMC
- Change to SIMC Makefile:

```
## CEBAF DEFAULT SETUP FLAGS:
simcdir = .

## THE REST SHOULD BE OK WITHOUT MODIFICATION.

## This tells make not to delete these target files on error/interrupt (see man page)
.PRECIOUS: *.o sos/*.o hms/*.o hrs1/*.o hrsr/*.o shms/*.o calo/*.o

RM      = rm -f
SHELL   = /bin/sh
S       = $(simcdir)/sos/
H       = $(simcdir)/hms/
L       = $(simcdir)/hrs1/
R       = $(simcdir)/hrsr/
A       = $(simcdir)/shared/
SH      = $(simcdir)/shms/
T       = $(simcdir)/cteq5/
C       = $(simcdir)/calo/
CH      = $(simcdir)/cern/
D       = $(simcdir)/fdss/

OBJ1    = target.o brem.o gauss1.o NtupleInit.o NtupleClose.o enerloss_new.o
OBJ2    = radc.o init.o dbase.o physics_kaon.o physics_pion.o physics_delta.o physics_proton.o loren.o sf_lookup.o
OBJ3    = semi_physics.o rho_physics.o rho_decay.o generate_rho.o trg_track.o semi_dilution.o
OBJ4    = results_write.o event.o call_ranlux.o jacobians.o F1F2IN21.o
OBJ5    = $(A)musc.o $(A)musc_ext.o $(A)project.o $(A)transp.o
OBJ6    = $(A)rotate_haxis.o $(A)rotate_vaxis.o $(A)locforunt.o
OBJ7    = $(H)mc_hms.o $(H)mc_hms_hut.o $(H)mc_hms_recon.o
OBJ8    = $(S)mc_sos.o $(S)mc_sos_hut.o $(S)mc_sos_recon.o
```

SIMC modifications

- Replace physics_pion.f with physics_iterate.f in SIMC

```
libra:/home/huberg/r2d2/simc/simc_fpi2
File Edit View Search Terminal Help
* Models for sigL, sigT, sigLT, sigTT for Deuterium.
***
* Parameterization revised for IT26, 12.11.09
  q2_set=2.45
  tav=(0.0735+0.028*log(q2_set))*q2_set
  ftav=(abs(t_gev)-tav)/tav
  ft=t_gev/(abs(t_gev)+0.139570**2)**2

  sigl=(fitpar(1)+fitpar(2)*log(Q2_g))
1      *exp((fitpar(3)+fitpar(4)*log(Q2_g))*(abs(t_gev)-0.2))
  sigt=fitpar(5)+fitpar(6)*log(Q2_g)
1      +(fitpar(7)+fitpar(8)*log(Q2_g))*ftav

  siglt=(fitpar(9)*exp(fitpar(10)*abs(t_gev))
1      +fitpar(11)/abs(t_gev))*sin(thetacm)
  sigtt=(fitpar(12)*Q2_g*exp(-Q2_g))*ft*sin(thetacm)**2

  tav=(-0.178+0.315*log(Q2_g))*Q2_g

  sig219=(sigt+main%epsilon*sigl+main%epsilon*cos(2.*phicm)*sigtt
>      +sqrt(2.0*main%epsilon*(1.+main%epsilon))*cos(phicm)*siglt)/1.d0

c now convert to different W
c W dependence given by 1/(W^2-M^2)^2
c factor 15.333 is value of (w**2-ami**2)**2 at W=2.19

c      wfactor=15.333/(s-(targ.Mtar_pion/1000.))**2)**2
*      wfactor=8.539/(s-(targ.Mtar_pion/1000.))**2)**2

c      wfactor=1.D0/(s-(targ.Mtar_pion/1000.))**2)**2
      wfactor=1.D0/(s_gev-mtar_gev**2)**2
      sig=sig219*wfactor
      sigl=sigl*wfactor
      sigt=sigt*wfactor
      sigtt=sigtt*wfactor
      siglt=siglt*wfactor
```

- **fitpar()** is an array of free parameters that will be determined as part of the L/T–iteration process.
- It is the functional dependence that needs to be determined first, not the parameters.
- However, it is also essential that you make some good guesses of initial parameter values for your first iteration.

Recalculating SIMC Weights

- You do not need to re-run SIMC every iteration if you have a way to over-write the weights
- Note, the weight is not $d^2\sigma/dtd\phi$!
- It is proportional to $d^5\sigma/dE'd\Omega e'd\Omega\pi$ and corresponds to the expected rate for the input luminosity!

See [event.f](#)

```
! The total contributing weight from this event -- this weight is  
! proportional to # experimental counts represented by the event.  
! Apply survival probability to kaons if we're not modeling decay.
```

```
main%weight = main%SF_weight*main%jacobian*main%gen_weight*main%sigcc
```

↑
Event mapping
Jacobian

↑
 $d^5\sigma$ from
physics routine

```
!  
! Note that there are also jacobians associated with some and/or all of  
! the above.  
! 1: We generate uniformly in xptar/yptar, not theta/phi. We define the  
! phase space volume (genvol contribution) as the product of the xptar/yptar  
! range, and have a jacobian for each event taking into account the mapping  
! between the solid angle on the unit sphere, and the dxptar/dyptar volume  
! (the jacobian is  $1/\cos^3(d\theta)$ , where  $d\theta$  is the angle between the  
! event and the central spectrometer vector)
```

Recalculating SIMC Weights

- You do not need to re-run SIMC every iteration if you have a way to over-write the weights
- Need to retain all factors going from $d^2\sigma/dtd\phi \rightarrow$ Weight when doing the re-calculation
- Hopefully you have access to one of Bill's Weight (ROOT/C++) re-calculation routines.
- This is my example from Fpi-2 LD2 analysis (PAW/Fortran):

```
lichen> more wt32.f
function wtn(q2_set)
include ?

real p1,p2,p3,p4,p5,p6,p7,p8,p9,p10,p11,p12
real nsigl,nsigt,nsiglt,nsigtt,tmp
real nsig219,nsig,wtn
real ft,tav,ftav

real pi,mtar_gev,q2_gev
real my_limit
integer q2_set
parameter (pi=3.14159)
parameter (mtar_gev=0.93827231)

if (abs(q2_set-245).lt.1) then
  p1= 0.28952E+02
  p2= -0.10000E+02
  p3= -0.15000E+02
  p4= 0.00000E+00
  p5= 0.46602E+02
  p6= -0.30000E+02
  p7= 0.18368E+01
  p8= 0.00000E+00
  p9= 0.10000E+04
  p10=-0.28000E+02
  p11= 0.35000E+01
  p12= 0.00000E+00
else
  write(*,*)'wtn: q2 error ',q2_set
endif
```

```
***
* Parameterization based upon Fpi-1 pi+ IT25, 12.04.18
* Revised for IT21, 12.11.06
*   tav=(0.0735+0.028*log(Q2i))*Q2i
  q2_gev=float(q2_set)/100.
  tav=(0.0735+0.028*log(q2_gev))*q2_gev
  ftav=(ti-tav)/tav
  ft=ti/(ti+0.139570**2)**2

*   nsigl=(p1+p2*log(Q2i))*exp((p3+p4*log(Q2i))*ti)
  nsigl=(p1+p2*log(Q2i))*exp((p3+p4*log(Q2i))*(ti-0.2))
  nsigt=p5+p6*log(Q2i)+(p7+p8*log(Q2i))*ftav

*   nsiglt=p9*exp(p10*ti)*sin(thetacmi)
  nsiglt=(p9*exp(p10*ti)+p11/ti)*sin(thetacmi)
  nsigtt=(p12*Q2i*exp(-Q2i))*ft*sin(thetacmi)**2

  nsig219=(nsigt+epsiloni*nsigl+epsiloni*cos(2.*phicmi)*nsigtt
1      +sqrt(2.0*epsiloni*(1.+epsiloni))*cos(phicmi)*nsiglt)/1.d0

  wfactor=1.D0/(Wcmi**2-mtar_gev**2)**2
  nsig=nsig219*wfactor

  nsig=nsig/2./pi/1.d+06 !dsig/dtdphicm in microbarns/MeV**2/rad

  wtn=Weight*nsig/dsigdt

  my_limit=0.20
  if ((wtn.lt.my_limit).and.(wtn.gt.0.0)) then
    continue
  else
    wtn=0.
  endif

  return
end
```

It is crucial to keep organized

- Each Q^2, W should be done separately.
- It is too much to expect the procedure to work globally, we only need to properly take into account the kinematic variation across a single diamond at a time.
- Keep each iteration in a different directory, e.g. Q^2_xx/IT_yy
- Keep ALL output. Don't throw anything away!
- Example fitpar() for $Q^2=2.45$ LD+ iteration #11

```
libra> cd it11/  
libra> ls  
par.pl_245  
libra> cat par.pl_245  
 0.82135E+03 0.12595E+03 1 4.1  
-0.41000E+03 0.00000E+00 2 4.1  
-0.24615E+02 0.12942E+01 3 4.1  
 0.11100E+02 0.00000E+00 4 4.1  
 0.35925E+02 0.18332E+01 5 0.4  
-0.18000E+02 0.00000E+00 6 0.4  
 0.27316E+02 0.62567E+01 7 0.4  
-0.31000E+02 0.00000E+00 8 0.4  
 0.00000E+00 0.00000E+00 9 47.4  
-0.20000E+02 0.00000E+00 10 47.4  
-0.12451E+03 0.98541E+01 11 1.2  
 0.00000E+00 0.00000E+00 12 1.2  
libra> 
```


Evaluating Model Dependence

- Blok et al.,
PRC 78 (2008) 045202

- This is the Model Dependence to our L/T/LT/TT, not the Model Dependence to any final Form Factor values, which must be evaluated separately!
- **DO THIS AT THE END, AFTER YOUR CROSS SECTIONS ARE FINALIZED!**

Since the extracted separated cross sections depend in principle on the cross section model, there is a “model” systematic uncertainty. This uncertainty was studied by extracting σ_L and σ_T with different cross section models. Since the longitudinal and transverse cross sections in the model reproduce the experimental values to within 10%, these two terms were independently increased and decreased by 10% in the model. With these changes, the extracted σ_L and σ_T varied by less than 0.5%. For evaluating the model uncertainty due to the interference terms σ_{LT} and σ_{TT} , these terms were independently increased or decreased by their respective uncertainties, obtained when fitting the four structure functions, and L/T separations were done with the modified models. The contribution to the uncertainty of σ_L and σ_T of these two terms is between 1% and 8% and depends strongly on t . The latter value (at the largest values of $-t$) is comparable to the contribution of uncorrelated uncertainties to σ_L and σ_T .

A few pointers

- Read over the Blok paper VERY CAREFULLY
- A single Q^2, W iteration should only take 1-2 hours
- I was able to do several iterations in a day, between teaching and other duties
- The `fitpar()` should converge in a few iterations to give $0.5 < R < 2$
- The main work is in getting R to be acceptably flat and in getting stable cross sections